

# Pattern formation and Turing instability

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## Topics:

- Pattern formation through “symmetry breaking” and “loss of stability”
- Activator-inhibitor systems with diffusion

Turing proposed a mechanism for growth and development of patterns (morphogenesis) in biological systems (embryonic development, et al). According to Turing

- Active genes stimulate production/activation of chemical agents (morphogenes)
- Chemical reactions alone is too “symmetric” for pattern generation
- But diffusion-driven instabilities create initial patterns and those can lead to further development

Basic mathematical questions:

Q1: can diffusion be *stabilizing factor* for reactive system ? (Answer: yes in 1D and some other cases)

Q2: can diffusion destabilize a reactive system?

Under what conditions?

What are the resulting patterns?

Examples of “symmetry breaking”, and patterns

- Physical systems: Phase transition “solid-fluid-gas”, fluid & gas have large continuous symmetries (all rotations, translations), solid – a rigid (discrete) crystalline symmetry.
- Mathematics: Hopf bifurcation from “stable equilibrium” (symmetric?) -> “limit cycle” (pattern?)

Turing analysis involves

- i) Symmetric (e.g. spatially uniform) equilibria
- ii) Bifurcations in various parameters, e.g. diffusivity, domain size
- iii) Activation/inhibition type reactions

## Stabilizing diffusion

For single reactant  $u(x, t)$  there is no Turing instability.

$$\begin{aligned} u_t &= Du_{xx} + f(u); \text{ on } \mathbf{R} \text{ or } [0, l] + \text{BC} \\ u_t &= D\Delta u + f(u); \text{ on } \mathbf{R}^n \text{ or } \Omega + \text{BC} \end{aligned}, \text{ or multi-D}$$

Indeed, a “uniform” ( $x$ -independent) ODE  $\dot{u} = f(u)$  has either stable  $m = f'(u^*) < 0$ , or unstable ( $m > 0$ ) equilibrium. The diffusion will maintain a stable one, and it can stabilize the unstable one.

To check it we take the linearized problem about  $u^*$ , for  $v(x, t) = u(x, t) - u^*$

$$\begin{aligned} v_t &= Dv_{xx} + mv; \\ m &= f'(u^*) \end{aligned} \tag{1}$$

## Solution on $\mathbf{R} / \mathbf{R}^n$

$$v(x, t) = \int_{-\infty}^{\infty} \frac{e^{m t - (x-\xi)^2/4Dt}}{\sqrt{4\pi Dt}} v_0(\xi) d\xi \tag{2}$$

For  $m > 0$ ,  $v(x, t)$  grows with  $t$

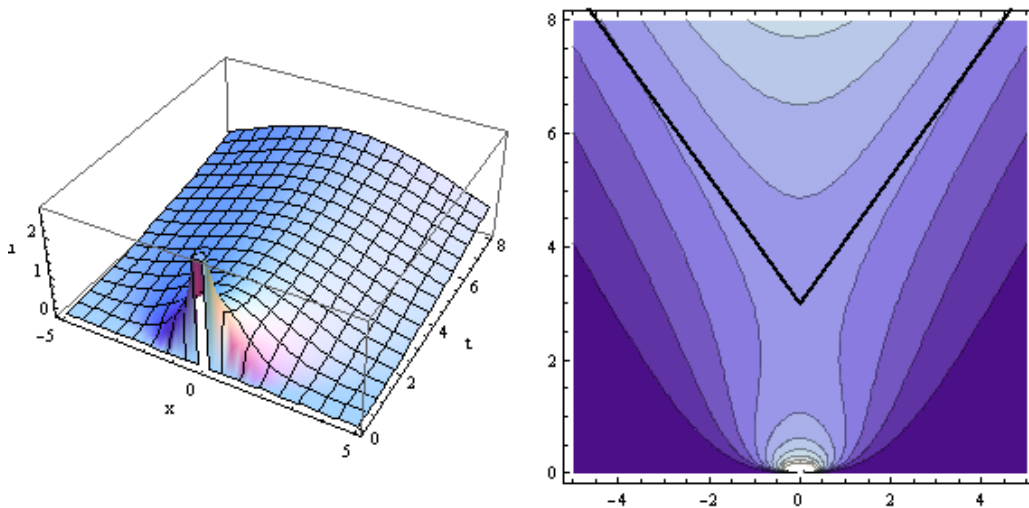


Figure 1: Point-source solution (2) on  $\mathbf{R}$  for  $m > 0$

For  $m < 0$ , all  $v(x, t)$  decay exponentially ( $v(x, t) < Ce^{-mt}$ ).

## Finite interval $[0, l]$

Use Neumann BC (no flux) and eigenfunction expansion

$$\begin{cases} \partial_x^2 \psi + \mu \psi = 0; \text{ on } [0, l] \\ \partial_x \psi|_{0, l} = 0 \end{cases} \Rightarrow \begin{cases} \mu_k = \left(\frac{\pi k}{l}\right)^2 \\ \psi_k(x) = \cos\left(\frac{\pi k x}{l}\right) \end{cases}; k = 0, 1, \dots \quad (3)$$

The eigenvalues of BV problem (1) are  $\mu_k = -D(\pi k / l)^2 + m$ . So stable case ( $m < 0$ ) remains stable, For unstable case get bifurcation values at  $D(\pi / l)^2 = m; D(2\pi / l)^2 = m; \dots$

The resulting would be “unstable mode” (pattern) is the corresponding cos-Fourier mode (3).

## Turing instability for double-diffusive systems

A pair of reaction-diffusion species obeys a coupled system

$$\begin{cases} u_t = D_1 u_{xx} + f(u, v) \\ v_t = D_2 v_{xx} + g(u, v) \end{cases} \quad (4)$$

or its multi-D version in  $(x, y, \dots)$  space, with Neumann BC. So symmetric (x-independent)  $(u, v)$  solve a pure reaction ODS with equilibrium  $(u^*, v^*)$ , and Jacobian

$$J = \begin{bmatrix} f_u^* & f_v^* \\ g_u^* & g_v^* \end{bmatrix}$$

Then we get a linearized system for  $u'(x, t) = u(x, t) - u^*; v'(x, t) = v(x, t) - v^*$  by

$$\begin{cases} u'_t = D_1 u'_{xx} + au' + bv' \\ v'_t = D_2 v'_{xx} + cu' + dv' \end{cases}; \text{ with Jacobian } J = \begin{bmatrix} a & b \\ c & d \end{bmatrix}$$

or in vector notation

$$U_t = \begin{bmatrix} D_1 & \\ & D_2 \end{bmatrix} \partial_x^2 U + J \cdot U; U(x, t) = \begin{pmatrix} u' \\ v' \end{pmatrix} \quad (5)$$

Stability of system (5) depends on eigenvalues of a *matrix-differential operator*

$$L = \begin{bmatrix} D_1 & \\ & D_2 \end{bmatrix} \partial_x^2 + J \quad (6)$$

We denote by  $\mu$  eigenvalues of scalar laplacian  $L_0 = \partial_x^2$  or  $\Delta$ .

**Examples of eigenvalue problem for Laplacian:**  $\Delta \psi + \mu \psi = 0 \rightarrow \{\mu, \psi(x)\} : ?$

	$\mu$	$\psi(x)$
Free space $\mathbf{R}; \mathbf{R}^n$	$0 \leq \mu_k = k^2 < \infty$ $\mathbf{k} = (k_1, k_2, \dots) \in \mathbf{R}^n;$ $k =  \mathbf{k} $	$\{\psi_{\mathbf{k}}(\mathbf{x}) = e^{i\mathbf{k} \cdot \mathbf{x}}\}$ or $\{\cos \mathbf{k} \cdot \mathbf{x}, \sin \mathbf{k} \cdot \mathbf{x}\}$
Finite interval $[0, l]$		
Dirichlet:	$\{\mu_k = (\pi k / l)^2 : k = 1, 2, \dots\}$	$\psi_k = \sin\left(\frac{\pi k x}{l}\right)$
Neuman:	$\{\mu_k = (\pi k / l)^2 : k = 0, 1, \dots\}$	$\psi_k = \cos\left(\frac{\pi k x}{l}\right)$
Periodic	$\{\mu_k = (2\pi k / l)^2 : -\infty < k < \infty\}$	$\{\psi_k(x) = e^{i2\pi k x / l}\}$ or $\{\cos, \sin\}$
2D (nD) Square, box $0 < x < a; 0 < y < b$	Product type: $\{\mu_{k,m} = \pi^2 [(k/a)^2 + (m/b)^2] : k, m = 1, 2, \dots\}$	$\psi_{km} = \sin\left(\frac{\pi k x}{a}\right) \sin\left(\frac{\pi m y}{b}\right)$
Polar disk (annulus) $0 < r < a;$ $0 < \theta < 2\pi$	product type (radial x angular modes): $\mu_{km} = (z_{km} / a)^2$	$\psi_{k,m}(r, \theta) = J_m(z_{km} r / a) e^{\pm i m \theta}$

Having compute eigenvalues of  $L_0$  we can pass to matrix-operator (6). We search for eigenmodes of  $L$  in the form "scalar x vector" =  $\psi(\mathbf{x})X$ ,  $\psi$  - eigenmode of Laplace  $L_0$ . It gives

$$\left( \lambda I + \mu \begin{bmatrix} D_1 & \\ & D_2 \end{bmatrix} - J \right) \cdot X = 0; \quad J = \begin{bmatrix} a & b \\ c & d \end{bmatrix}$$

The eigenvalue problem for "diagonalized" operator  $L$  is then reduced to 2D matrix-problem for a family of matrices

$$B(\mu, J) = J - \mu \begin{bmatrix} D_1 & \\ & D_2 \end{bmatrix} = \begin{bmatrix} a - \mu D_1 & b \\ c & d - \mu D_2 \end{bmatrix} \quad (7)$$

Specifically, eigenvalues of  $L$

$$\rightarrow \lambda(L) = \lambda(\mu, J) = \text{eigen}[B(\mu, J)] \quad (8)$$

where  $\{\mu\}$  vary over the entire spectrum of Laplacian.

### Turing instability analysis

Basic question: given positive Laplace eigens  $\{\mu\}$  can a stable chemistry (Jacobian  $J$ ) produce an unstable matrix  $B(\mu)$  for specific diffusivities  $D_{1,2}$ , domain size  $l$ , or "norm" of  $J$  (proportional to reaction rates in  $f, g$ ).

Table 1: stability test

Stable J	Stable B- ?
$\text{tr}J = a + d < 0$	$\text{tr}B = -\mu(D_1 + D_2) + \text{tr}J < 0$ - always true!
$\det J = ad - bc > 0$	$\det B = D_1 D_2 \mu^2 - \underbrace{(aD_2 + dD_1)}_{a_1} \mu + \underbrace{(ad - bc)}_{a_2} = Q(\mu)$

The answer depends on quadratic function  $Q(\mu)$  -? Specifically,

**Case 1:** Having symmetric (negative-definite)  $J = \begin{bmatrix} a & b \\ b & d \end{bmatrix}$ , matrix function  $B(\mu)$  (7) is negative-

definite for all  $\mu > 0$ . So  $\lambda(\mu, J)$  remain negative, and diffusion leads to stabilization (no patterns).

**Case 2:** Scalar diffusivity  $D_1 = D_2 = D$ . Clearly, for stable  $J$  and  $\mu > 0$ , "eigens  $B$ " = "eigens( $J$ ) -  $\mu D$ " are more negative (stable).

In general case (Turing) it depends on coefficients of quadratic function  $Q(\mu) = \det B(\mu)$  in Table 1. Namely, (i)  $Q(\mu) > 0$  for all  $\mu$  means "stable  $B(\mu)$ ", hence stable L; (ii)  $Q(\mu)$  changing sign for some  $\mu > 0$ , gives unstable  $B(\mu)$ , hence "unstable mode" for L. Turing conditions for "unstable" (negative)  $p(\mu)$  are

$$\begin{aligned} (T1) \quad & a_1 = aD_2 + dD_1 > 0 && \text{- negative slope } Q'(0) \\ (T2) \quad & (aD_2 + dD_1)^2 \geq 4D_1D_2(ad - bc) && \text{- discriminant} \end{aligned} \quad (9)$$

From (T1) and  $\text{tr}(J) < 0 \Rightarrow$  coefficients (a,d) should have opposite signs, then (b,c) are also opposite. The only choice are

**Case 3:** Activation-inhibition systems: (i) "u" grows, "v" decays; (ii) "u" inhibits "v" and "v" enhances "u", or vice versa

$$J = \begin{bmatrix} + & - \\ + & - \end{bmatrix}; \text{ or } \begin{bmatrix} + & + \\ - & - \end{bmatrix}$$

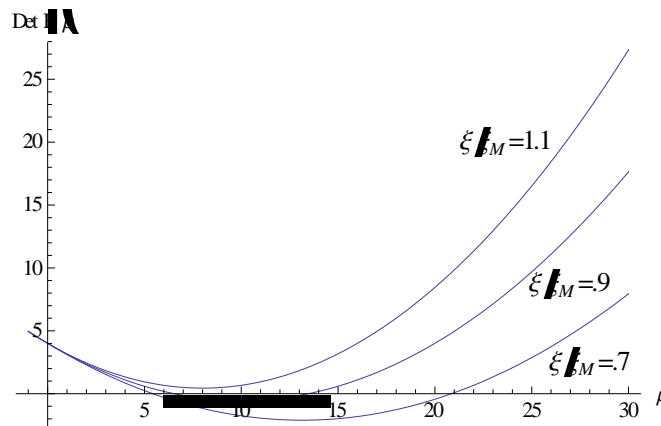


Figure 2: Cases of stable and unstable  $\det(B)$  in terms of parameter  $\xi = \sqrt{D_1 / D_2}$

One can rewrite Turing condition (9) for  $J = \begin{bmatrix} a & \pm b \\ \mp c & -d \end{bmatrix}$  (a,b,c,d – positive), as

$$a / \xi + d\xi \geq 2\sqrt{bc}; \text{ with parameter } \xi = \sqrt{D_1 / D_2} \quad (10)$$

Combined set conditions ( $\text{tr}(J) < 0, \det(J) > 0, \text{tr}(B) < 0, \det(B) < 0$ ) gives

$$a / d \leq 1;$$

$$\xi = \sqrt{D_1 / D_2} \leq \frac{\sqrt{bc} - \sqrt{bc - ad}}{d}$$

It sets a limit (Fig.3) of ratio  $D_1 / D_2$  to get unstable Turing mode.

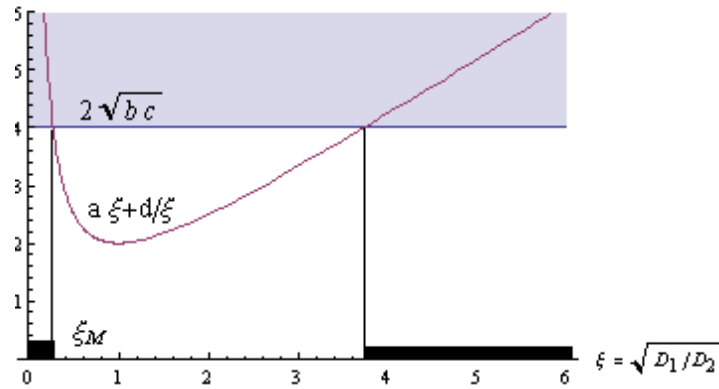


Figure 3: Parameter range for  $D_1/D_2$  to get Turing instability

A positive (unstable) eigenvalue  $\lambda(L)$  for some  $\mu$  within the “unstable” interval of Fig.2 will give a spatial (eigenmode) part  $\psi(\mathbf{x})$  (e.g.  $\cos\left(\frac{\pi kx}{l}\right)$  for  $= -\partial_x^2$ , with  $\mu_k = \left(\frac{\pi k}{l}\right)^2$ ) to “seed” the follow up development of a **spatial pattern** for nonlinear reaction-diffusion (4), spawn by its “symmetric” unstable equilibrium  $(u^*, v^*)$ .

**Bifurcation analysis** in terms of  $D_i$ , depends on how “parabola”  $\sigma(\mu)$  of Fig.2 will hit the positive  $\mu$  axis. It can be “saddle-node” or Hopf bifurcation(?).