

*STOCHASTIC MODELS IN ENGINEERING AND SCIENCE*  
*October 10-11, 2008*

# **Geostrophic turbulence: How Jupiter got its stripes?**

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# Applied Math projects at Case in 90s

- Center for Stochastic and Chaotic Processes
  - ONR project on "Stochastic modeling in oceanography"
  - CRDF projects on "Geophysical modeling"
    - V. Klyatskin
    - S. Danilov**
    - F. Dolzhansky
- IAP, Moscow, Russia

## Topics:

- Methods of Statistical topography in Turbulent transport
1. **V. Klyatskin, D. Gurarie, W. Woyczynski**, *Diffusing passive tracers in random incompressible flows: statistical topographic aspects*, **J. Stat. Phys.**, 1996, 84, #3/4, 797-836.
  2. **D. Gurarie, V. Klyatskin, W. Woyczynski**, *Short-time correlation approximations for diffusing tracers in random velocity fields: a functional approach*, *Stochastic modeling in physical oceanography*, **Progress. in Prob.**, 39, Birkhauser, Boston, Boston, MA, 1996, 221-269, 1997
  3. V. Klyatskin, D. Gurarie, *Coherent phenomena in stochastic dynamical systems*, *Physics - Uspekhi Fizicheskikh Nauk (Russian Physics Surveys)*, 169, no.2, 1999, 171-207
  4. ....

# Lecture Plan

## **Geophysical flows**

rotation and gravity

Vorticity conservation and Quasigeostrophy

## **Turbulence**

3D and 2D phenomenology (direct and inverse energy cascades)

Rhines theory

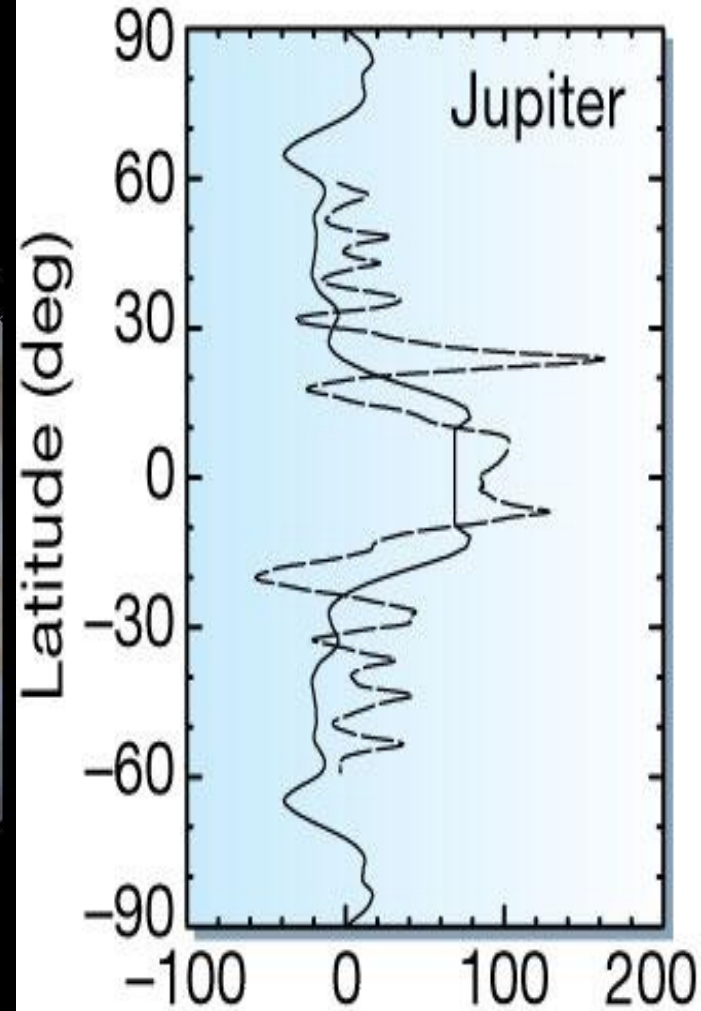
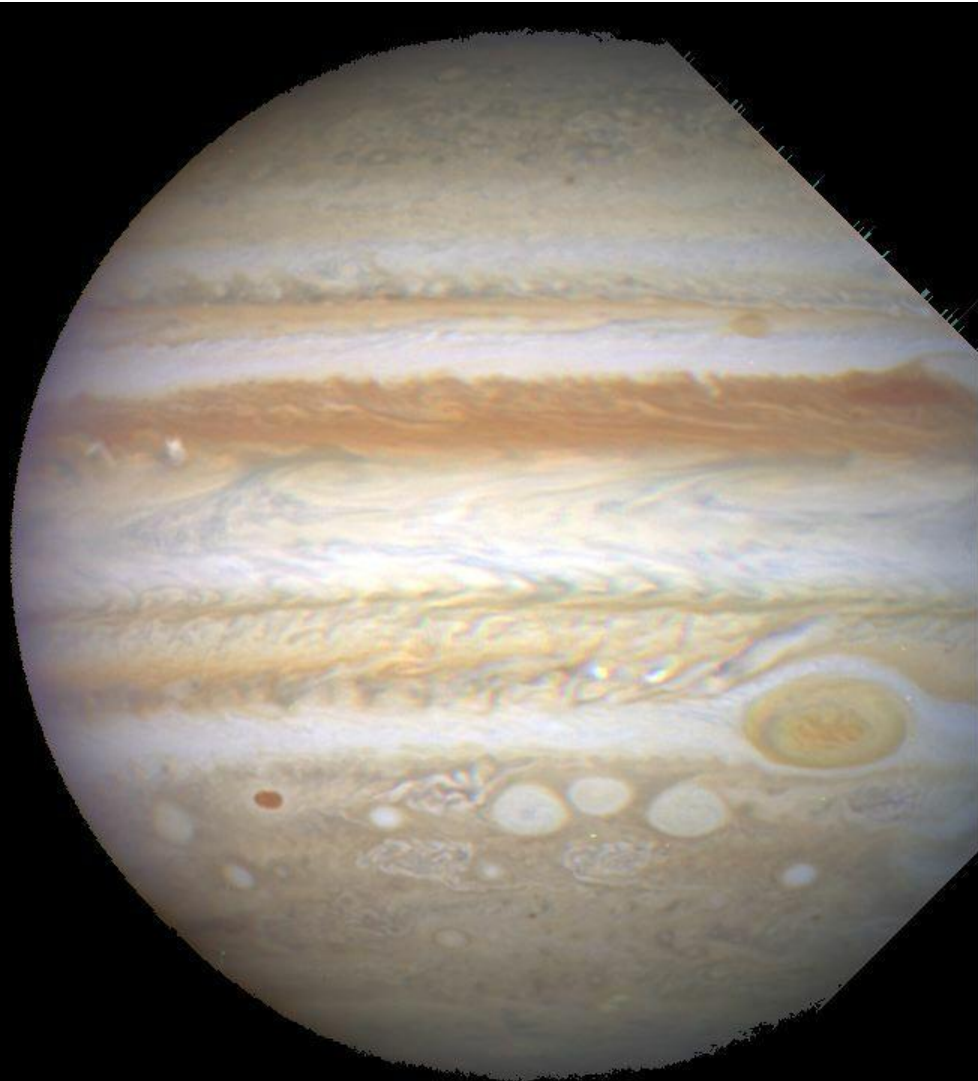
## **CFD “turbulence”**

Spectra and large-scale structure (jets and vortices)

Evolution of zonal jets

Open problems and issues

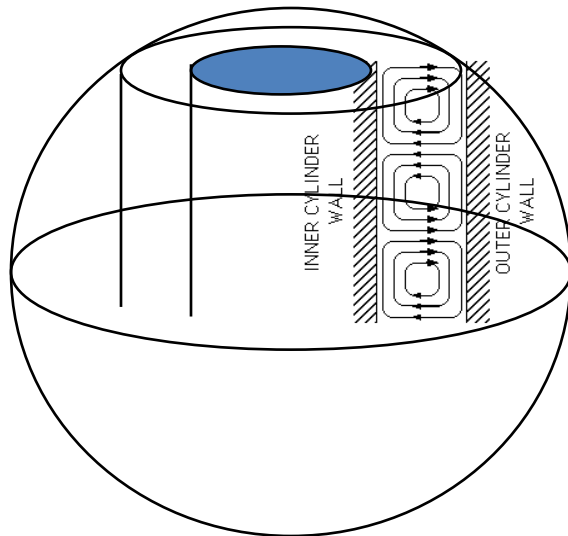
# Zonal jets on Jupiter



Zonal winds m/s  
Dashed - measured  
Solid - simulated

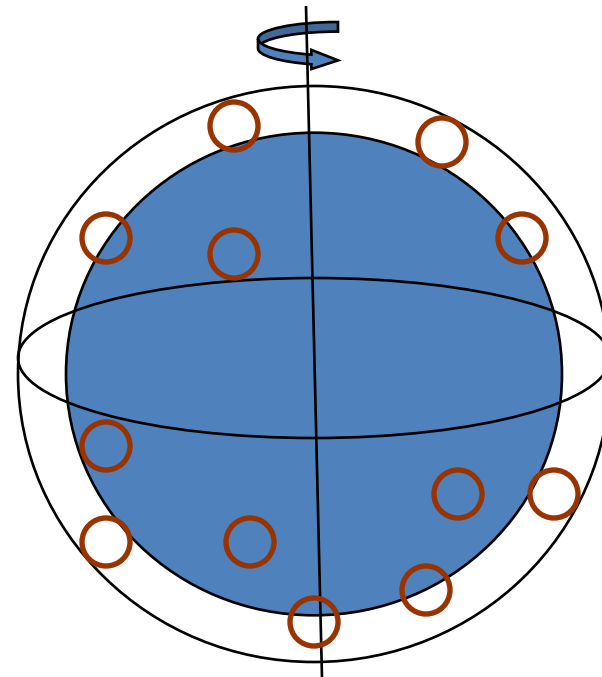
# Origin, mechanism ??

Taylor-Couette vortices in differentially rotating spherical (cylindrical) shells (Busse, 1983)



Deep convection

Quasigeostrophic 2D-turbulence, stirred by "small-scale" stochastic forcing (Rhines 75, Wilson 78, ...)



Shallow convection

# I. Basic fluid equations

Velocity:  $\vec{v} = (u, v)$  (2 D)  $(u, v, w)$  (3 D)

Stream:  $\psi(x, y) \rightarrow \vec{v} = (-\psi_y, \psi_x)$  (2 D)

Vorticity:  $\zeta = v_x - u_y = \nabla^2 \psi$

Momentum conservation:  $\partial_t \vec{v} + \vec{v} \cdot \nabla \vec{v} = -\frac{\nabla p}{\rho} + \nu \nabla^2 \vec{v} + \vec{F}$

Incompressibility / mass :  $\nabla \cdot \vec{v} = 0$

Energy cons.  $\dots$

$p$  - pressure;  $\nu$  - viscosity;  $\vec{F}$  - body forces

Vorticity form :

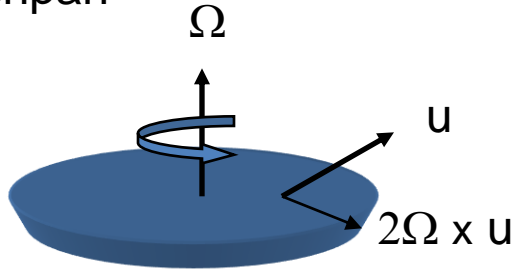
$$\partial_t \zeta + J(\psi, \zeta) = 0 \text{ or } \nu \nabla^2 \zeta,$$

$$J(f, g) = \det \begin{pmatrix} f_x & f_y \\ g_x & g_y \end{pmatrix} - \text{Jacobian}$$

Solutions ? Analysis ?

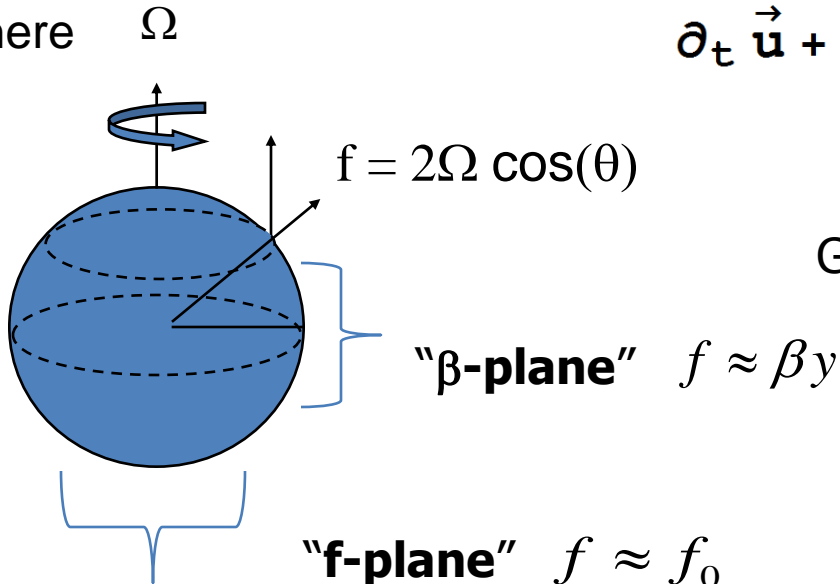
# Rotating fluids: Geostrophic balance

Dishpan



$f = 2\Omega$  – local Coriolis

Sphere



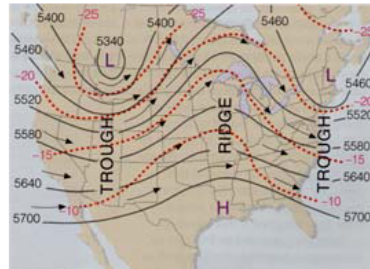
$$\partial_t \vec{u} + \vec{u} \cdot \nabla \vec{u} + \vec{f} \times \vec{u} = - \frac{\nabla p}{\rho} + \nu \nabla^2 \vec{u}$$

Geostrophic winds follow isobars

# Geostrophic winds follow isobars

## Pressure patterns and winds aloft

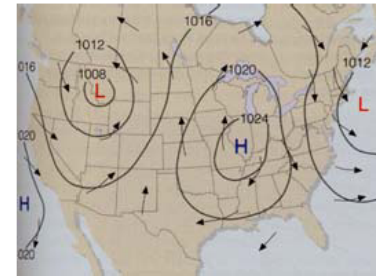
- At upper levels, winds blow parallel to the pressure/height contours



## Surface pressure patterns and winds

Near the surface in the northern hemisphere winds blow

- counterclockwise around and in toward the center of low pressure areas
- clockwise around and outward from the center of high pressure areas



Why doesn't the wind blow from high to low pressure?

Dominant factors/forces for large scale geophysical flows

- Rotation (Coriolis, geostrophic balance)
- Gravity (stratification, hydrostatic pressure)

# Word *geostrophic*

made of two Greek roots:

**Geo** (*earth*) + **Strophe** ( part of an ancient Greek choral ode sung by the chorus when moving from right to left, hence movement performed by the chorus during the singing of this part) – *cyclonic motion*

Ancient Greeks were interested in meteorology, and several examples of ***meteorological terms*** with Greek roots survived, though not their ***theories*** (Aristotle thought that winds from the west are cold because “*they blow from the sunset*”)

# Quasi-geostrophy (QGS)

- QGS = departure from Geostrophic balance
- **Method:** rescaling/expansion in small *Rossby number*  $Ro = U/(fL)$  (similar to “Mach” in compressible gas) for linearized PV about *geostrophic balance*, dropping “fast” (gravity) modes
- Rossby on Jupiter:  $.02 < Ro < .3$

# Conservation of Potential vorticity: derivation of quasigeostrophy

**Primitive** (mass-momentum) equations conserve *Potential Vorticity*

$$PV = \text{“planetary } f\text{”} + \text{“relative”} \quad \zeta = \nabla \times \mathbf{v}$$

**Example:**

Rotating Shallow  
Water on f-plane

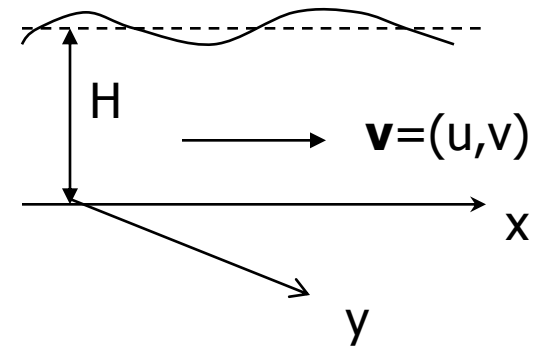
$$D_t \mathbf{v} + \mathbf{f} \times \mathbf{v} = -g \nabla h + \dots$$

$$h_t + \nabla \cdot [(H + h) \mathbf{v}] = 0$$

$g$  – gravity acceleration = 9.8m/s<sup>2</sup>

$H$  – mean height/depth

$h$  – departure from mean



$$PV: q = \frac{f + \zeta}{g(H + h)} \approx \frac{f}{gH} (1 + \zeta / f - h / H)$$

- **Quasigeostrophy** – small departure from *geostrophic balance*
- Expansion in small Rossby number  $Ro$  ( $.02 < Ro < .3$  for Jupiter!) for *linearized PV*

# Examples of QGS

$$q = \frac{f + \xi}{g h} \quad (\text{shallow water PV or its linearization})$$

$$\partial_t q + J(\psi, q) = D q + F$$

2 D (Euler) fluid :

$$\xi = \nabla^2 \psi$$

$$\partial_t \xi + J(\psi, \xi) = D[\xi] + F$$

Rotating dish pan "f-plane" :

$$q = (\nabla^2 - 1 / L^2) \psi$$

$$\partial_t q + J(\psi, q) = D[q] + F$$

Rotating sphere of radius R :

$$Q = \nabla^2 \psi + R \cos(\gamma / R)$$

$$\partial_t Q + J(\psi, Q) = D[Q] + F$$

Equatorial channel "beta-plane" :

$$q = \nabla^2 \psi + \beta y$$

$$\partial_t \nabla^2 \psi + J(\psi, \nabla^2 \psi) + \beta \partial_x \psi = D \psi + F$$

F – Forcing;

$$D = \lambda + \nu \nabla^2 - \text{"friction"}$$

# II. Turbulence

## Onset: shear instability

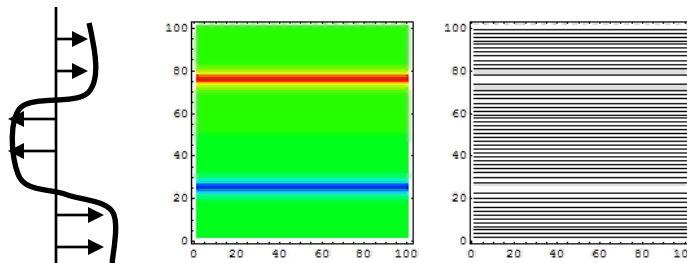
Shear profile :  $U(y) = -\Psi_y$ ;  $Z(y) = -U_y$ ;

Perturbation :  $\psi = \Psi(y) + \psi'$ ;  $\zeta = Z(y) + \zeta'$ ;

Linearized problem :  $\partial_t \zeta' + L[\zeta'] = 0$

Linear Operator :  $L = [ (\beta + U_{yy}) \nabla^{-2} - U ] \partial_x - D$

Kelvin-Helmholtz rolls



Planetary vortex (beta term) creates effective "shear":  $Z(y) = \beta y$

# Kolmogorov-Obukhov theory (1941)

- **Scales and Mode** : fluid flow viewed as system of interacting (Fourier) modes
- **Energy cascade**: energy distributed among modes (spectrum  $E_k$  ), and exchanged (energy flux  $\varepsilon$ )

## Premises:

- Viscosity is important only at small scales, hence large/intermediate ones (*inertial range*) are “inviscid” (no energy loss)
- Mode interaction cascades energy from larger to smaller scales - *direct cascade*
- In the absence of “distorting forces” energy spectra  $E_k$  should depend on scalar wave-number  $k$  (*isotropic spectrum*)
- **Locality**: only nearby scales (modes) exchange energy, hence constant (scale independent) *energy flux*  $\varepsilon$
- **Scaling laws** for  $E_k$  ?

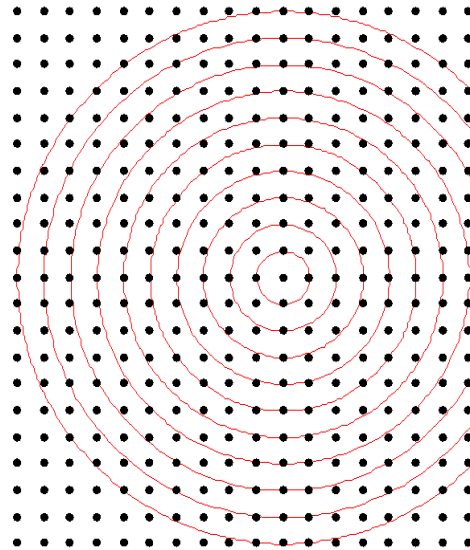
# Mode decomposition

PDEs (Euler, Navier-Stokes et al) converted to system of interacting modes

$$\psi(\mathbf{x}) = \sum_{\mathbf{k}=(\mathbf{k},l)} \psi_{\mathbf{k}} e^{i\mathbf{x}\cdot\mathbf{k}} - \text{Fourier modes (scale} = |\mathbf{k}|)$$

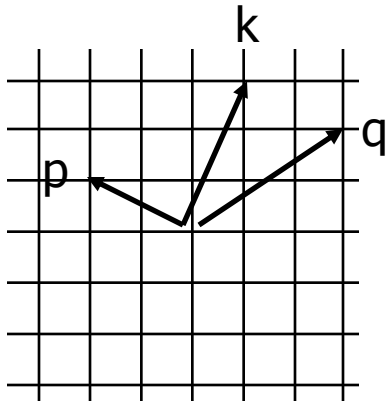
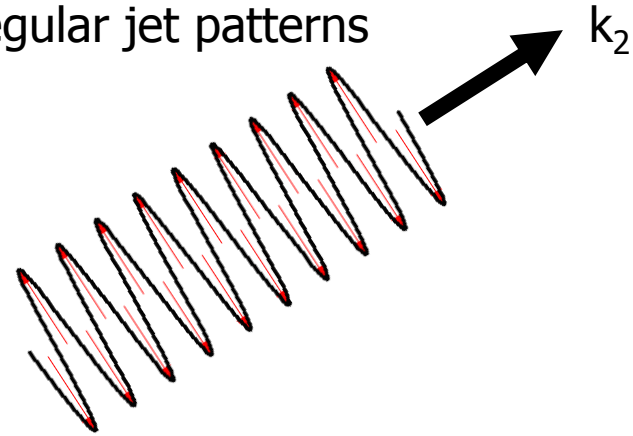
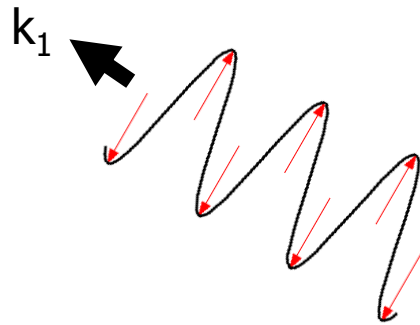
$$\text{2 D Euler : } \dot{\psi}_{\mathbf{k}} + D_{\mathbf{k}} \psi_{\mathbf{k}} = \sum_{\mathbf{p}+\mathbf{q}=\mathbf{k}} A_{\mathbf{p}\mathbf{q}} \psi_{\mathbf{p}} \psi_{\mathbf{q}} + F_{\mathbf{k}}$$

$$\text{Coefficients : } A_{\mathbf{p}\mathbf{q}} = \frac{\mathbf{p} \wedge \mathbf{q} (p^2 - q^2)}{k^2} - \text{interaction, } D_{\mathbf{k}} = \lambda + \nu k^2 - \text{dissipation}$$



Spectral energy shells (2D)

Fourier modes = regular jet patterns



PDE converted to coupled ODES for mode coefficients:

$$a_{km}(t)\cos(kx+my)+ \dots$$

# Dimensional argument [K41]

Energy flux:  $[\varepsilon] = \frac{m^2}{s^3}$ ; Energy density:  $[E_k] = \frac{m^3}{s^2}$ ; Wave number:  $[k] = \frac{1}{m}$

Scaling law:  $E_k \sim C\varepsilon^p k^q \Rightarrow p = \frac{2}{3}; q = -\frac{5}{3}$

$$E_k = C_K \varepsilon^{2/3} k^{-5/3}$$

## Development

- Theoretical approaches:

moment closure scheme (quasi-normal, etc.), RNG, stochastic models

- Fluid experiments and observation

- Numeric simulations

- Applications of turbulence/scaling laws:

engineering, atmospheric turbulence, transport, diffusion etc.

# 2D turbulence phenomenology:

## BKL (Batchelor-Kraichnan-Leith)

2D fluid (Euler) has many conserved integrals: “energy” + “moments of vorticity”

$$E = \frac{1}{2} \iint \mathbf{v}^2 dA = \sum_{\mathbf{k}} |\hat{\mathbf{v}}_{\mathbf{k}}|^2 - \text{energy}$$

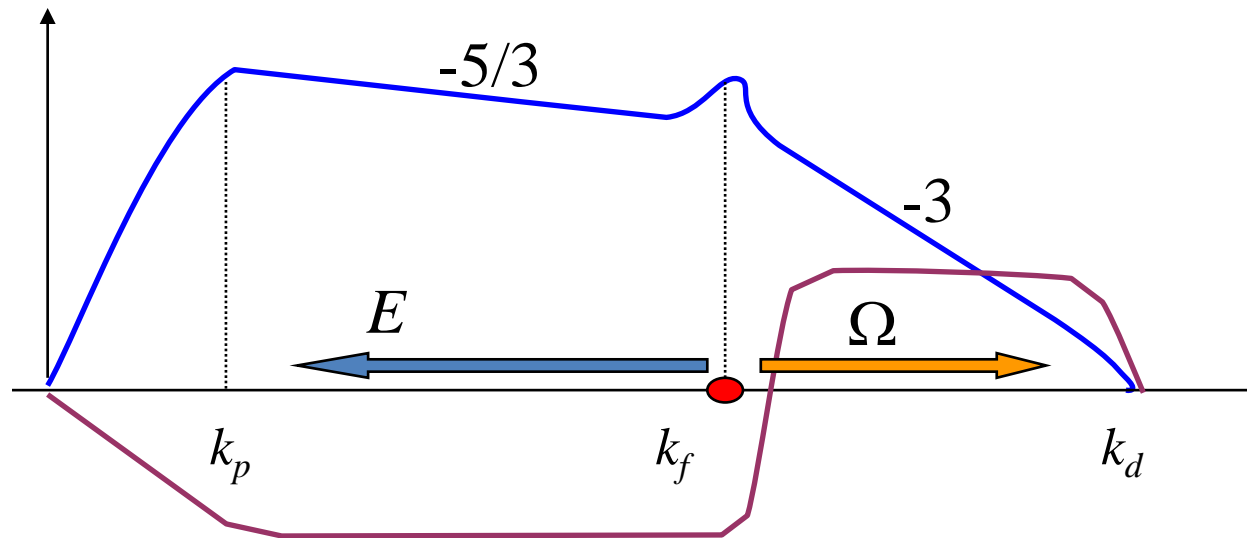
$$Z = \frac{1}{2} \iint \zeta^2 dA = \sum_{\mathbf{k}} k^2 |\hat{\mathbf{v}}_{\mathbf{k}}|^2 - \text{enstrophy}$$

=> Two (possible) inertial ranges and cascades:

**Direct cascade:** “large” to “small” scales (enstrophy conservation)

**Inverse cascade:** “small” to “large scales” (energy conservation)

# Isotropic Energy spectrum in 2D with source at $k_f$



## Direct cascade:

*"Big whorls have little whorls,  
Which feed on their velocity,  
And little whorls have lesser  
whorls,  
And so on to viscosity."*

(L. Richardson)

## Inverse Cascade:

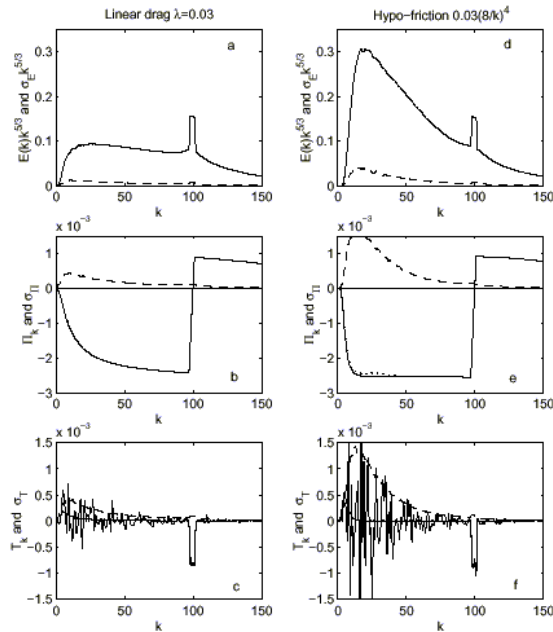
*"Big whorls meet bigger ones  
And so it tends to go on  
By merging they grow bigger yet  
And bigger yet, and so on..."*

(M.E. McIntyre)

**Inverse cascade leads to "large scale structures" in 2D flows !**

# Computational turbulence: Departure from BKL phenomenology

Numeric simulations on 512x512  
spectral grid



➤ Steep spectra (slope  $> 5/3$ )

➤ Nonuniform fluxes

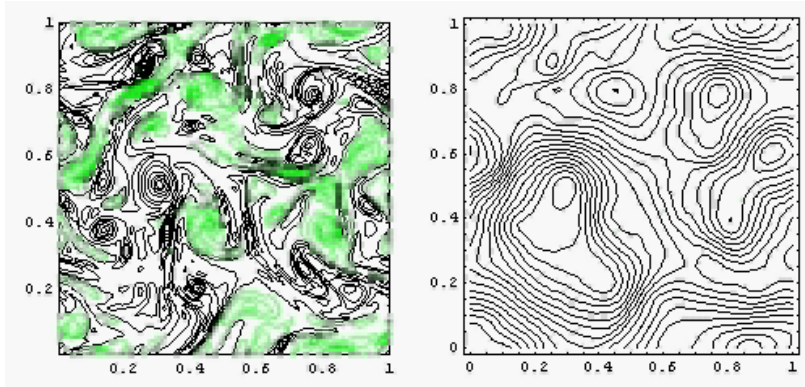
➤ Nonlocal transfers

- Vallis-Maltrud, 1991, 1993
- Borue, 1994
- Sukhoryanski et al 1999
- Boffetta et al, 2000
- Danilov-Gurarie, 2001-2002

# Computational turbulence: vorticity and stream field

I. Semi-Lagrangian code (DG) on 150x150 grid

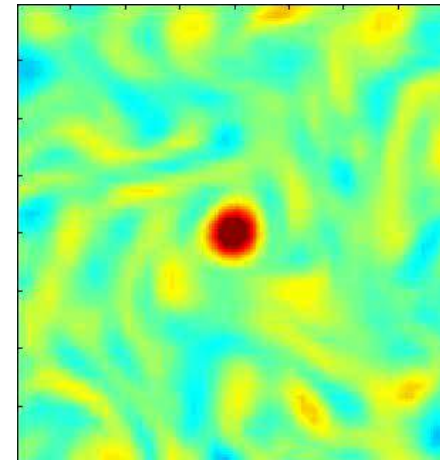
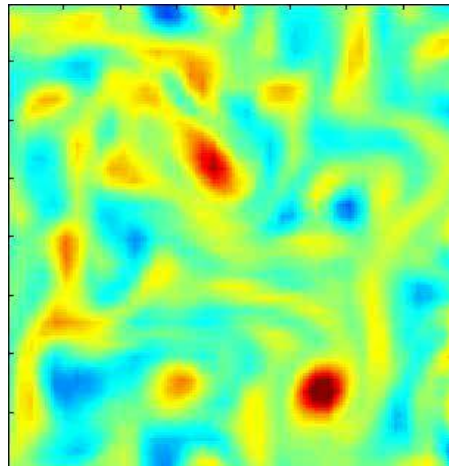
*(spatial discretization, and vorticity conservation)*



Stochastically forced turbulence (at  $k_f=80$ )

Strong vortex in forced turbulence

II. Spectral code (SD) on 512x512 Fourier grid  
*Fast FFT*



- ***Inverse cascade*** leads to **vortex mergers**
- Can it lead to **strong jets** ?

# III. Rhines theory of “beta-plane” turbulence

Beta-plane QGS combines “linear dispersion operator” + “nonlinear HD”

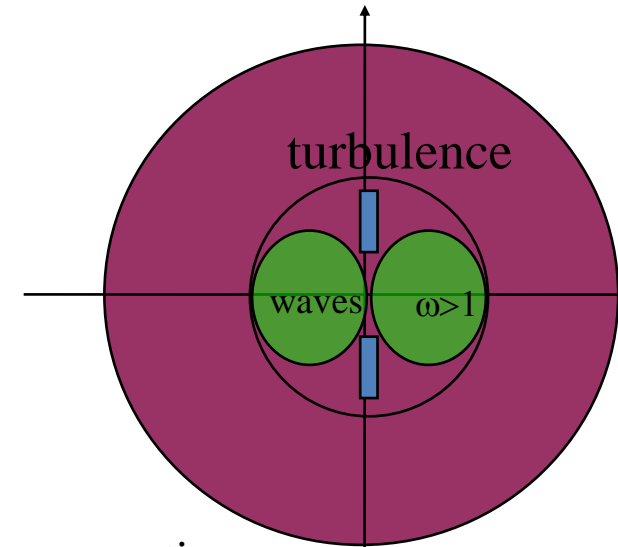
$$\partial_t \nabla^2 \psi + \underbrace{J(\psi, \nabla^2 \psi)}_{\text{NL term}} + \underbrace{\beta \partial_x \psi}_{\text{Linear}} = D \psi + F$$

Imposes additional constraint on triad interactions (“weak turbulence”)

$$\partial_t \psi_{\hat{k}} + i \omega_{\hat{k}} \psi_{\hat{k}} + \sum_{\hat{p}+\hat{q}=\hat{k}} A_{\hat{p}\hat{q}} \psi_{\hat{p}} \psi_{\hat{q}} = D_{\hat{k}} + F_{\hat{k}};$$

$$\omega_{\hat{k}} = \frac{\beta k_1}{k_1^2 + k_2^2} - \text{Rossby dispersion (waves)}$$

$\hat{k} = \hat{p} + \hat{q}$ - triad interaction $\omega_{\hat{k}} = \omega_{\hat{p}} + \omega_{\hat{q}}$ - wave resonances
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- “Beta” inhibits inverse cascade at scale through *Rossby wave generation*
- Suppresses cascade into high Rossby region (“lazy eight”) via “frequency detuning”
- Channels energy into zonal (near zonal) modes

# Frequency-detuning mechanism

(Holloway-Hendershot, 1977)

- Cascade into lazy-eight imposes frequency synchronism
- Triad transfers are dampened by Rossby-wave dispersion

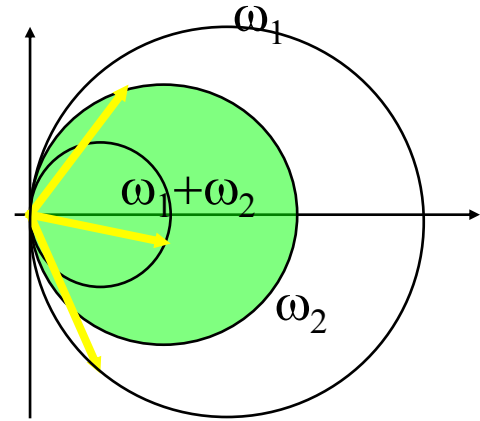
Rhines phenomenology:

“Eddy time-oscillations” = “*Rossby frequency*”

coupled with “Eddy speed” = “*Rossby phase speed*”

yield “*Rhines scale*” = “arrest scale” for inverse cascade on beta-plane

$$k_{\text{Rh}} = \sqrt{\frac{\beta}{2 U_{\text{rms}}}} = \frac{\beta^{1/2}}{2 E^{1/4}}$$

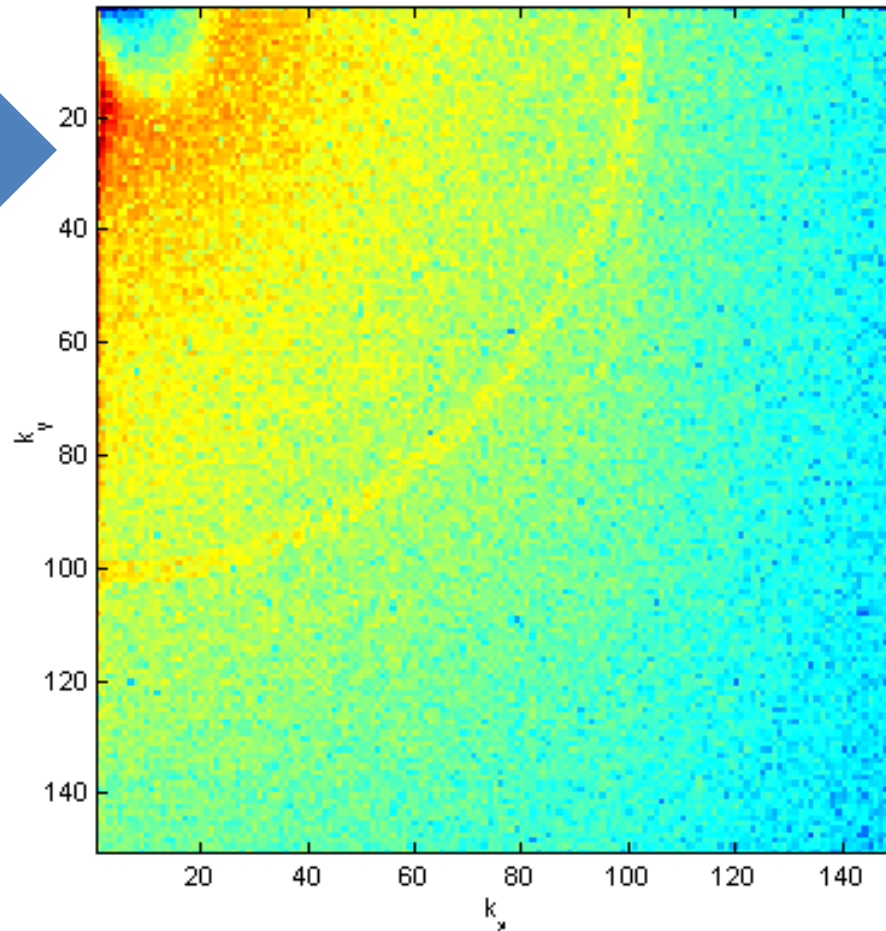
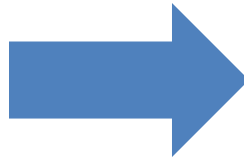


Also Rhines predicted  $k^{-5}$  zonal spectrum at large scales (?)

# Computational work: DNS on beta-plane (...; DG)

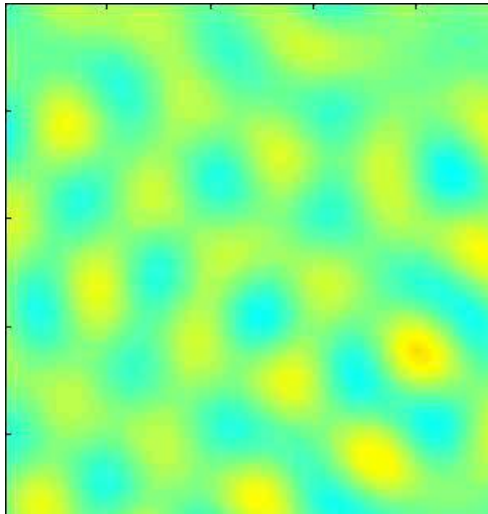
Beta-plane spectrum

**Spectral anisotropy** - accumulation of zonal modes

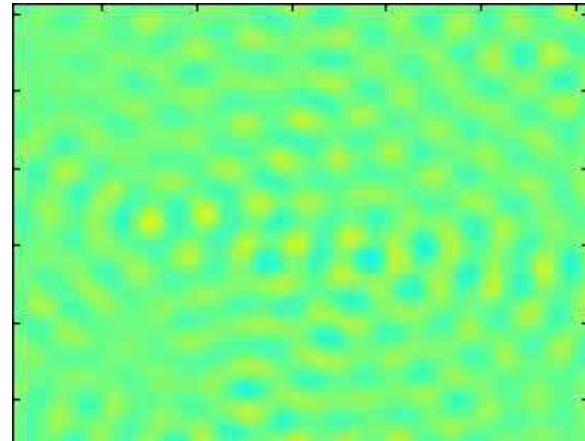


# Zonal jets: initial growth phase

Isotropic case

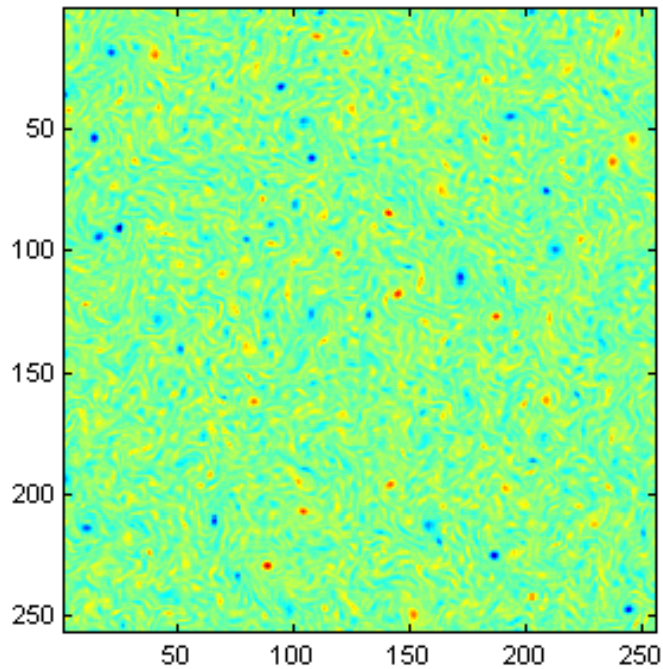


Beta- plane



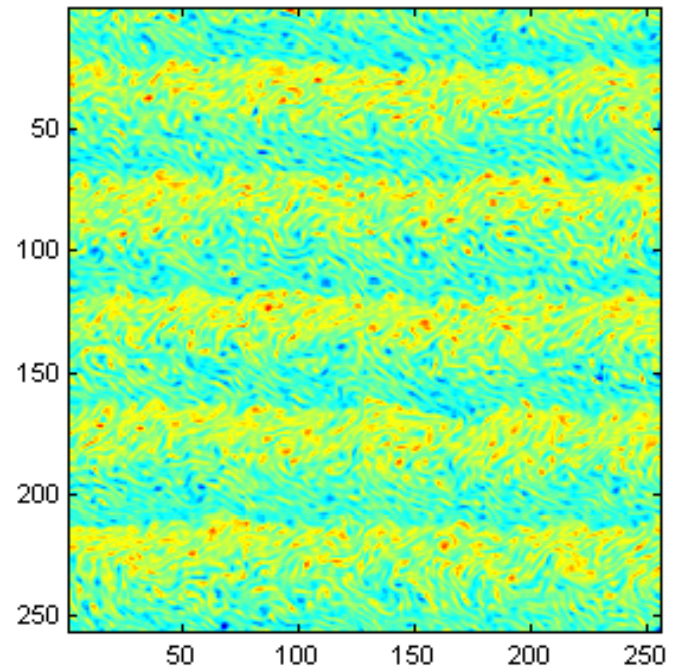
# Typical vorticity fields (DG)

Isotropic ( $\beta=0$ )



Strong localized vortices

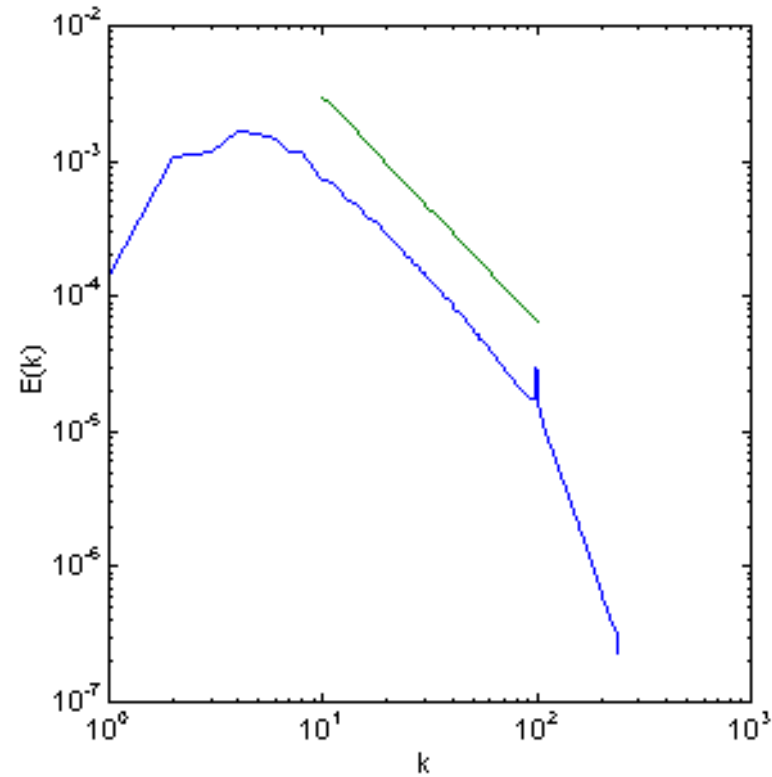
Beta



Zonal jets

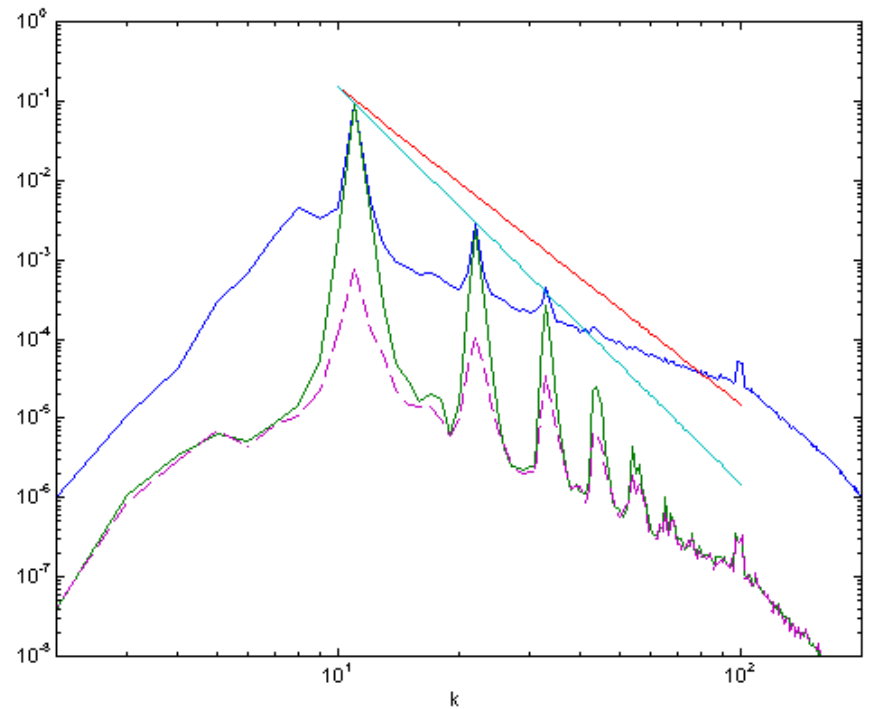
# Typical energy spectra (DG)

Isotropic



$5/3$  slope

Beta

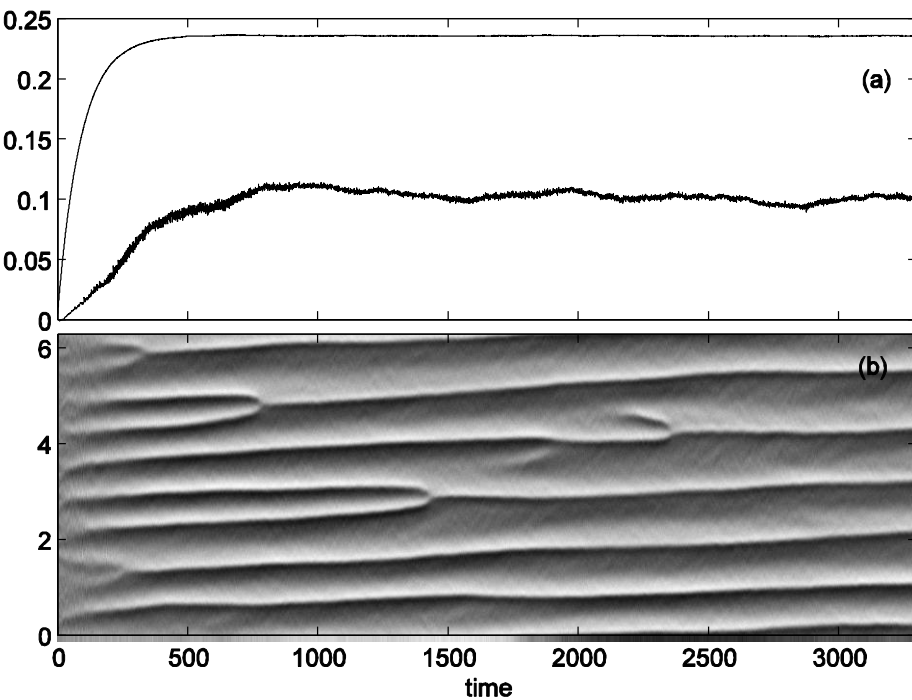


$-4$  and  $-5$  slopes

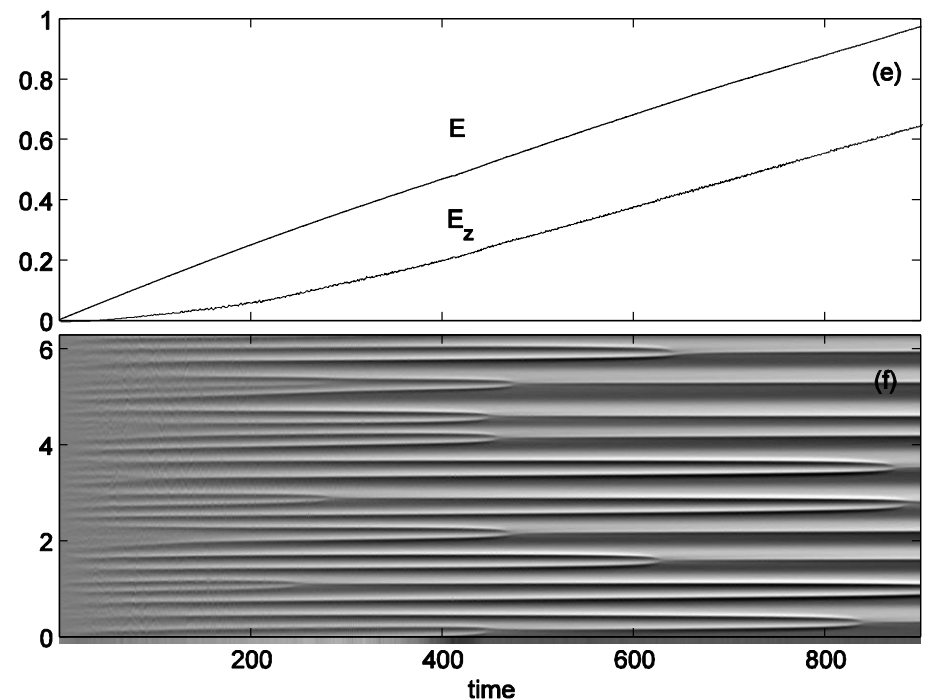
Dominated by zonal spectral peaks

# Long term evolution of zonal jets

Typical friction case



Frictionless case

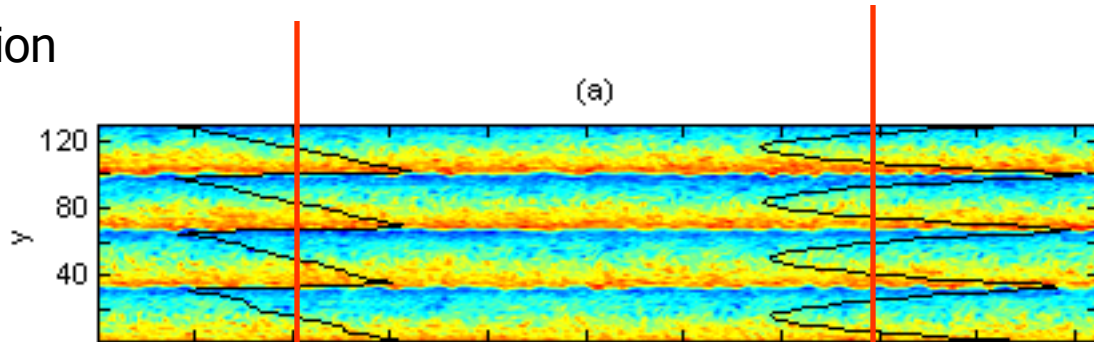


Slow emergence and evolution through jet mergers  
Stabilizes at  $k_{Rh}$

# Vorticity field and Zonal profiles (Jovian jets ?)

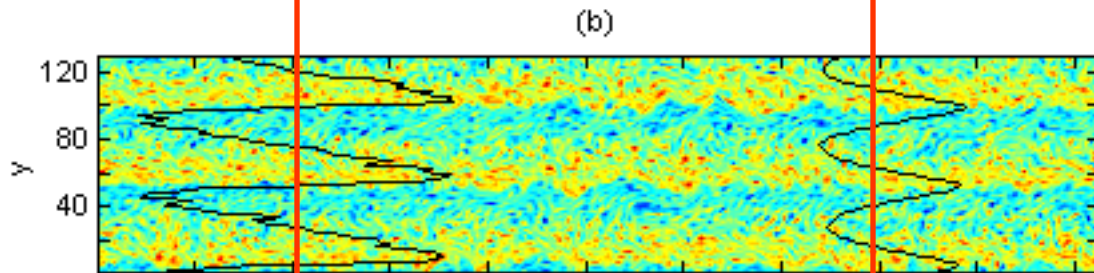
Frictional dissipation

Hypofriction  
 $\beta = 320$



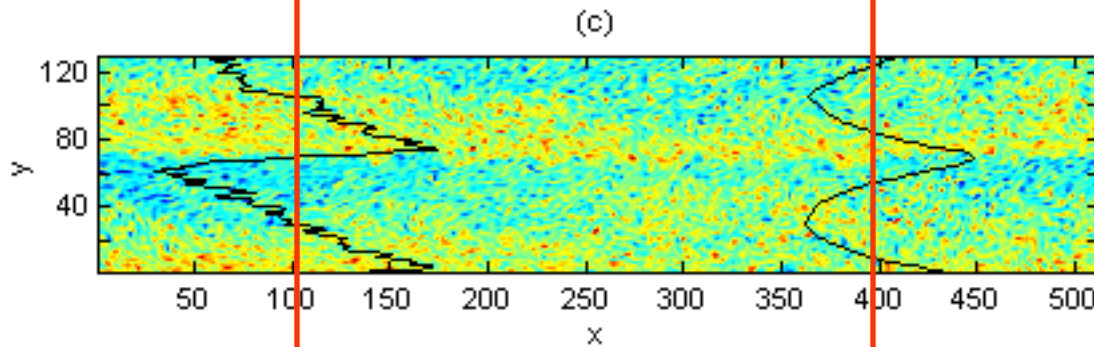
“slope”/beta  
= .6

Hypofriction  
 $\beta = 80$



.5

Standard  
(bottom) friction  
 $\beta = 80$

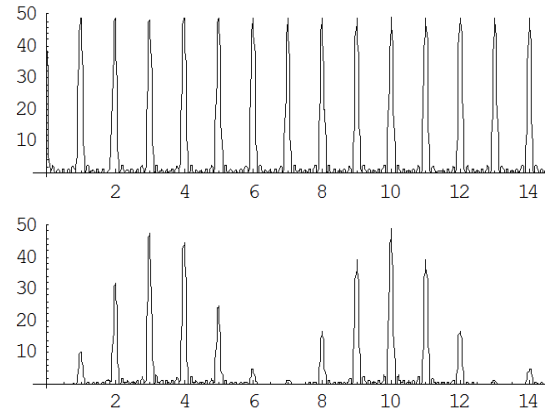
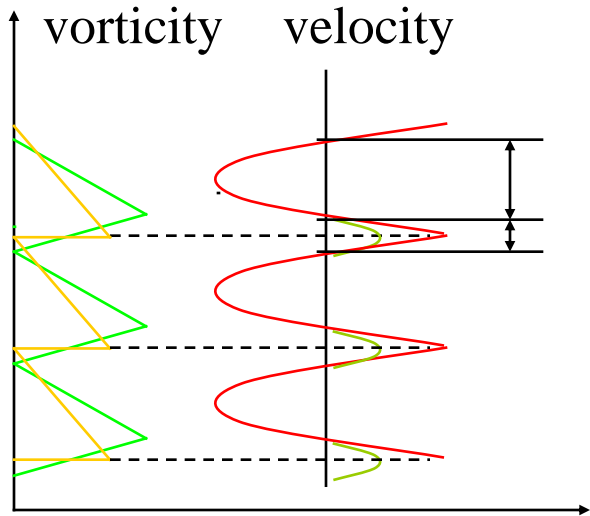


.2

Vorticity  $\zeta=0$

$U=0$

# Idealized zonal profiles and “scaling laws”



$$W_{\text{East}}/W_{\text{west}} = 2.3$$

Model shape functions:

$$\Phi_N(x) = \begin{cases} |\sum_{m=0}^N e^{2\pi i m x}|^2; & \text{saw} \\ |\sum_{m=0}^N (e^{2\pi i m (x-a)} - e^{2\pi i m x})|^2. & \text{slopi} \end{cases}$$

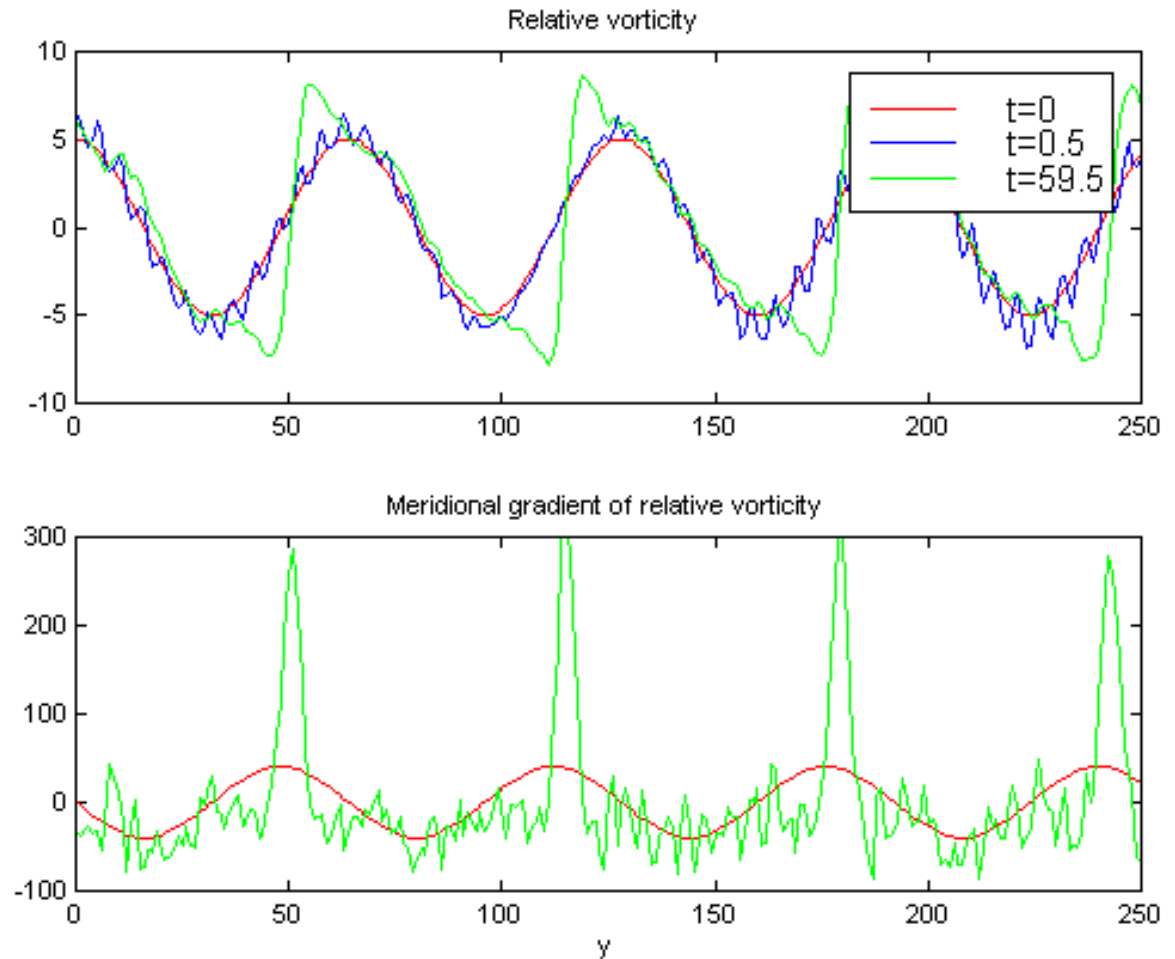
Zonal spectra (based on “shape function”)

$$E_{\text{zon}}(k) = \frac{C \alpha^2 \Phi(k/k_{\text{Rh}})}{k^p k_{\text{Rh}}^{-q}}$$

- “**Scaling laws**” for zonal spectra (peaks) have sensitive dependence on “jet patterns”. Could vary between  $k^{-4} - k^{-6}$ .
- Proposed “ **$k^{-5}$  law**” is special case (nothing “universal”)

# “Stare-case” zonal attractor

Initial symmetric profile  $\zeta = \sin(k y)$  with  $k = k_{Rh}$  develops “stare-case” pattern



# Theory of slow zonal evolution / equilibration:

Manfroy-Young (1999), Stuhne (2001)

## **Premises/method:**

- “Beta effect” on transverse (meridional) jet in weakly “supercritical regime” (past “linearized stability”)
- Expansion in “small parameter” = supercriticality (above bifurcation level)
- “Zonal jet equation” arises as slow (modulation) amplitude DE of Cahn-Hilliard type.

Interesting ideas, but ...

- Premises do not apply to strong turbulent regime with forcing/dissipation.
- Zonal jets don't grow as “*instabilities*” of large scale regular (meridional) flow, but arise as “*instabilities*” of small-scale turbulent “5/3 background”
- The proper theory/model - still wanting!

# Conclusions

*"... the turbulence community has been facing crisis through emergence of **coherent structures** in numerical and laboratory modelings so that the standard theories based on **homogeneity** and **isotropy** are increasingly challenged..."*

(Quote from GTP/NCAR workshop on "Cumulus Parameterization in the Context of Turbulence", 2004)

- Stochastically forced 2D turbulence (Euler/ QGS) channels energy to large scales (Inverse Cascade).
- The large-scale coherent structures on beta-plane is array of zonal jets, at Rhines scale.
- Slow zonal evolution leads to a "stare-case" vorticity profile of (negative) slope:  $-\alpha > -\beta$ , and the corresponding "strong" eastward and "weak" westward jets.
- The "stare-case" pattern serves as an "attractor" of stochastically forced turbulent flows.
- Mathematical challenge: explain (develop model of) "zonal instability" of "5/3 turbulent background" that gives rise to jet structure, its slow evolution and persistence

Are Jupiter stripes “zonal jets” of QGS turbulence

(“Busse” vs. “QGS”)

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